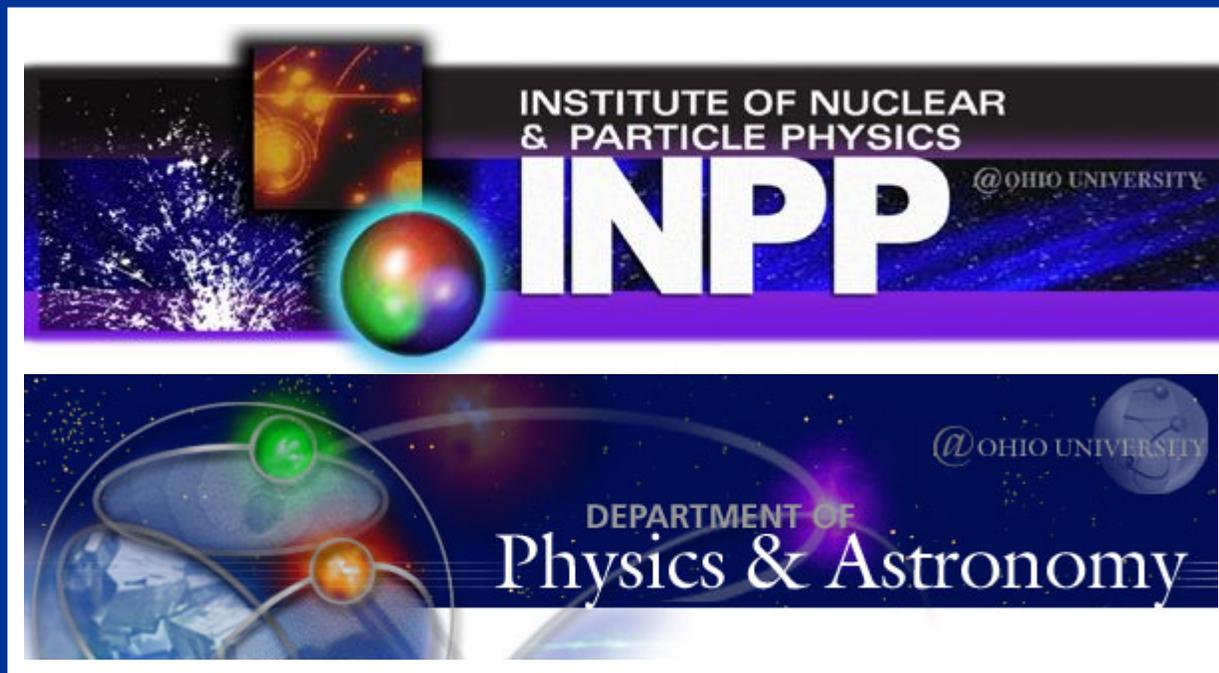


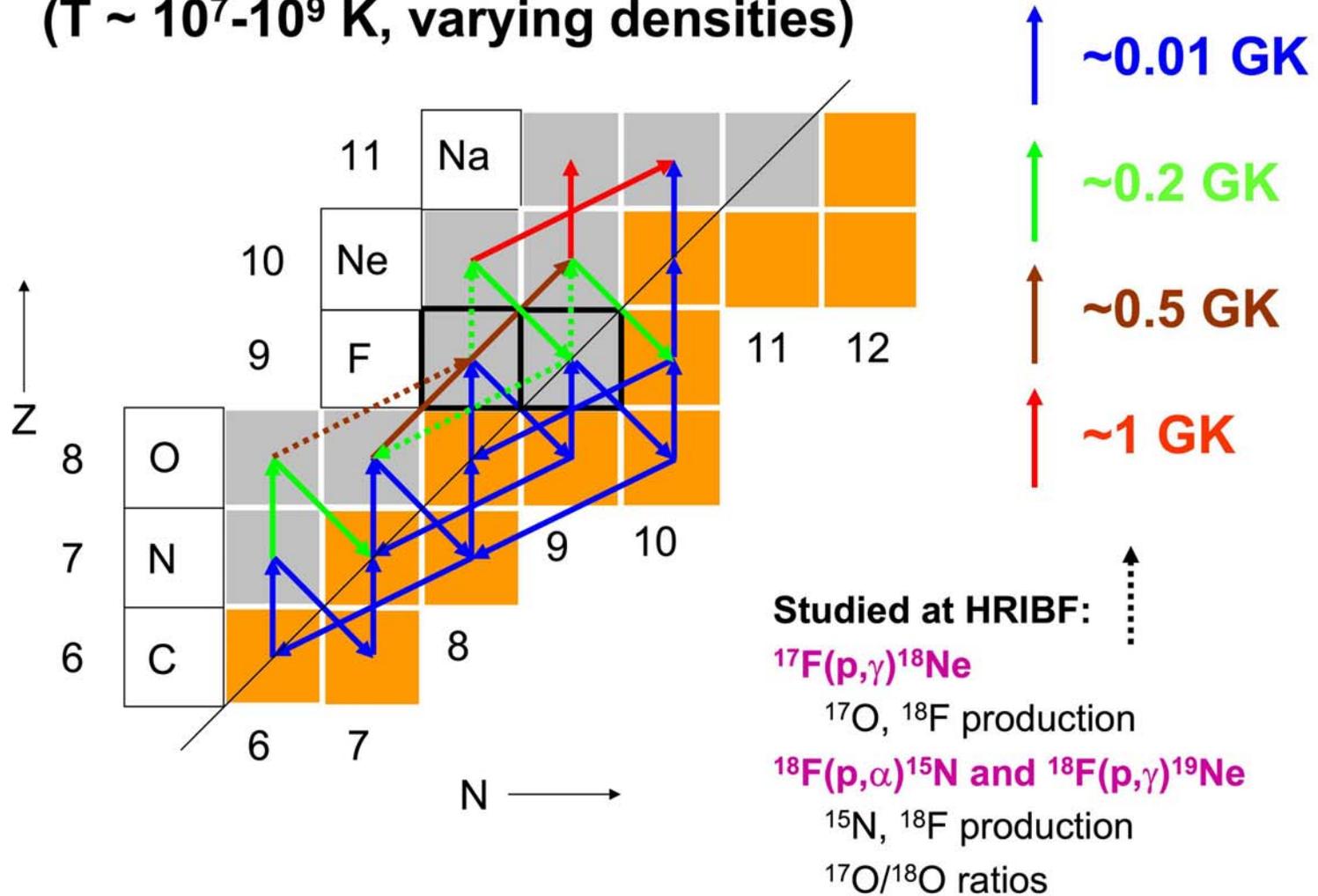
Nuclear Astrophysics - II

Carl Brune



Reactions in Stellar Explosions

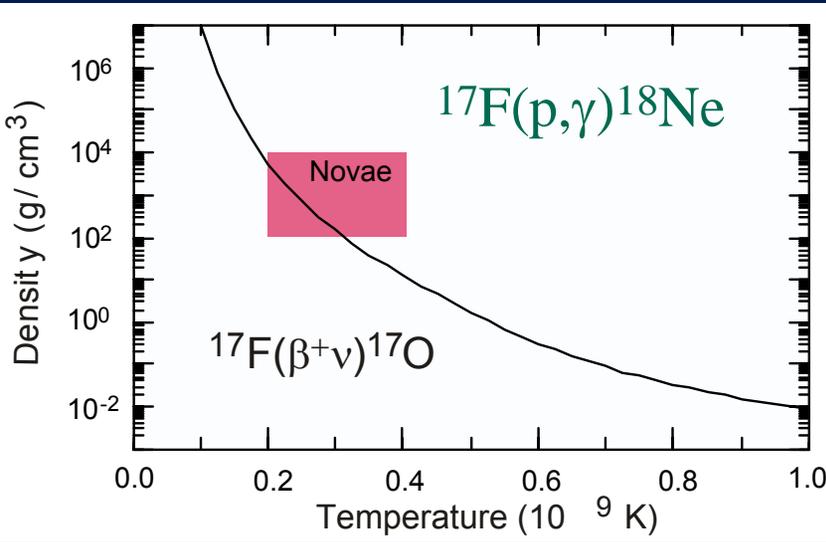
($T \sim 10^7$ - 10^9 K, varying densities)



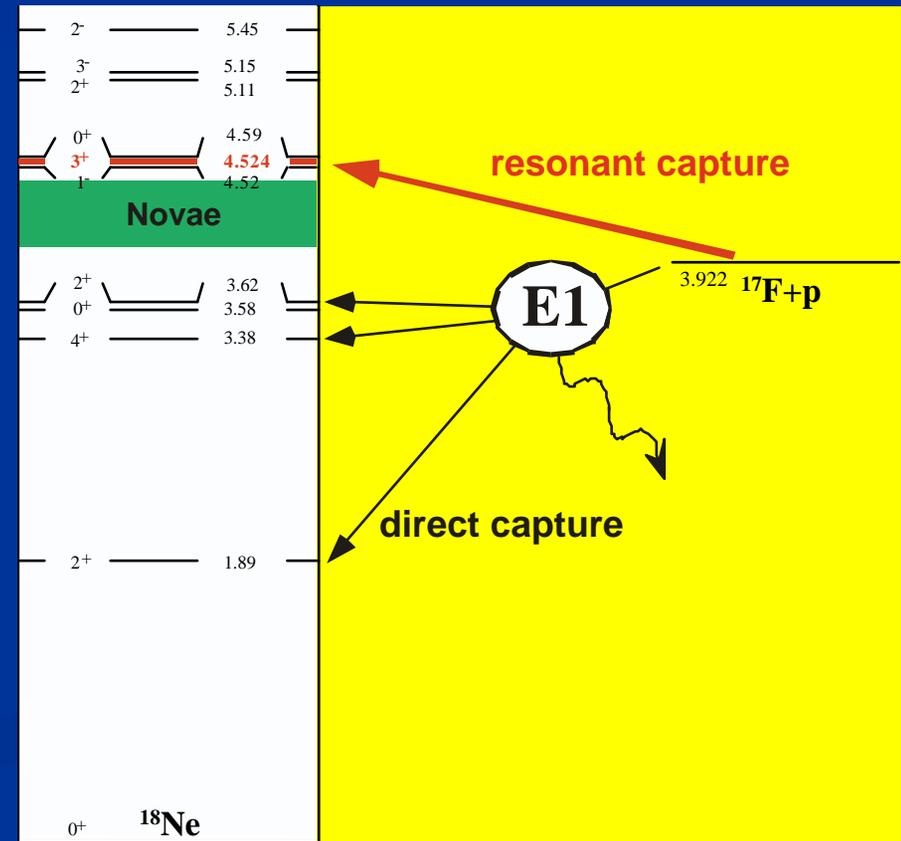
The $^{17}\text{F}(p,\gamma)^{18}\text{Ne}$ reaction

rate

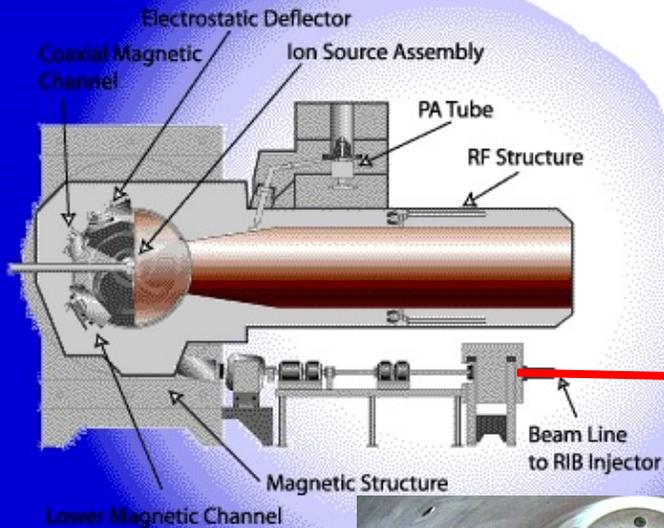
- Capture rate is comparable to the beta decay rate in novae.



- Two contributions to the rate:
 - Direct capture
 - 3+ resonance

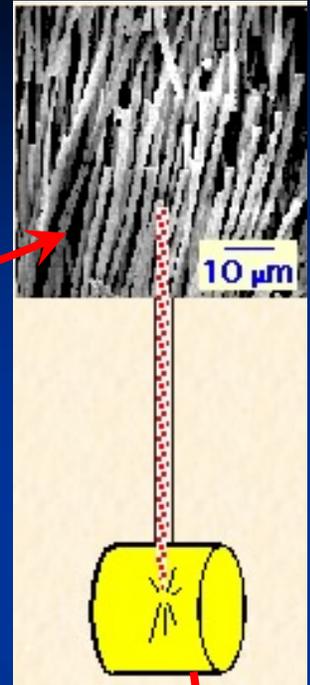


The Holifield RIB Facility at Oak Ridge National Lab



p, d, or a

Hot, fibrous production target



Ion source

25 MV tandem

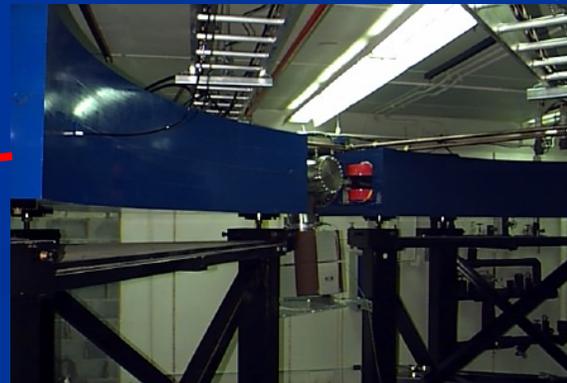
Mass analysis

RIB (300 keV)



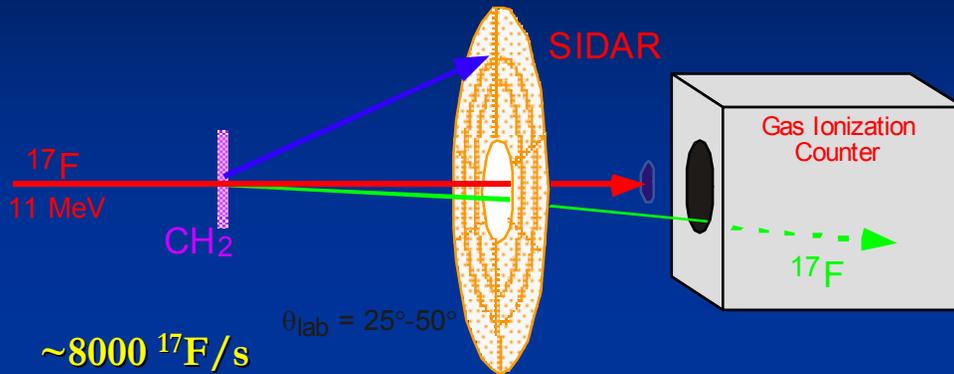
ORIC

To experiments

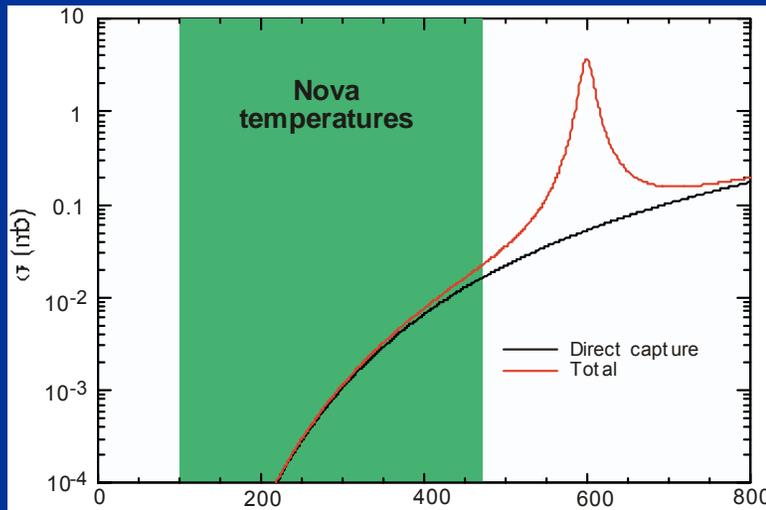
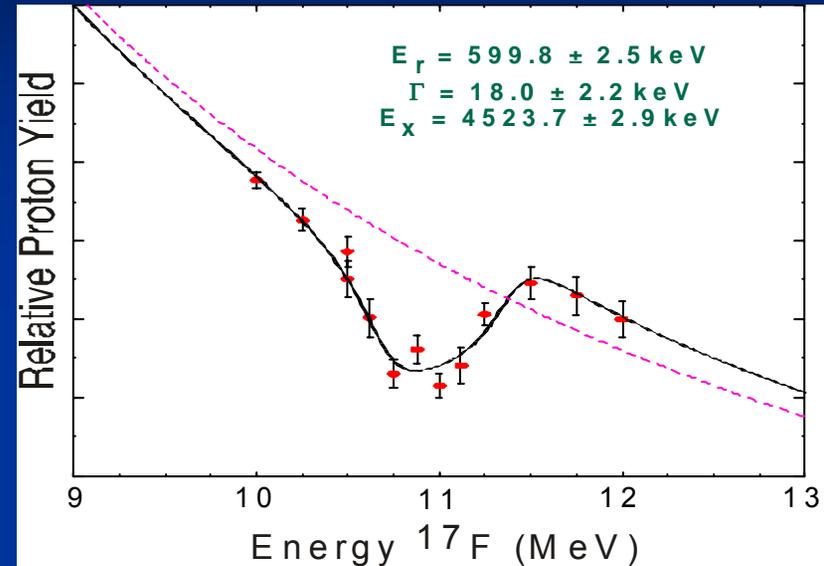


Energy and width of 3^+ measured via $^{17}\text{F} + \text{p}$ scattering

D. W. Bardayan *et al.*, Phys. Rev. C **62** (2000) 055804.
D. W. Bardayan *et al.*, Phys. Rev. Lett. **83** (1999) 45.



- ~ 8000 $^{17}\text{F}/\text{s}$
- $48 \mu\text{g}/\text{cm}^2$ polypropylene (CH_2) target
- Protons detected in large Silicon Detector ARray (SIDAR)
- Heavy ions detected in coincidence by ionization counter



- 3^+ resonance is too high in energy to contribute significantly to the rate at nova temperatures.
- Direct capture dominates, but cross section is unmeasured. Estimates based on ^{18}O .
- Direct capture cross section is too small to be measured at available ^{17}F intensities.

THE $^{17}\text{F}(p, \gamma)^{18}\text{Ne}$ DIRECT CAPTURE CROSS SECTION

J. C. Blackmon^a, D. W. Bardayan^a, C. R. Brune^b, F. C. Carstoiu^c,
A. E. Champagne^d, R. Crespo^e, T. Davinson^f, J. C. Fernandes^e,
C. A. Gagliardi^c, U. Greife^g, C. J. Gross^a, P. A. Hausladen^a, C. Iliadis^d,
C. C. Jewett^g, R. L. Kozub^h, T. A. Lewis^a, J. F. Liang^a, B. H. Moazen^h,
A. M. Mukhamedzhanov^c, C. D. Nesaraja^h, F. M. Nunesⁱ, P. D. Parker^j,
D. C. Radford^a, L. Sahin^d, J. P. Scott^h, D. Shapira^a, M. S. Smith^a,
J. S. Thomas^k, L. Trache^c, R. E. Tribble^c, P. J. Woods^f, and C.-H. Yu^a

^a *Oak Ridge National Laboratory, Oak Ridge, TN, USA*

^b *Ohio University, Athens, OH, USA*

^c *Texas A&M University, College Station, TX, USA*

^d *University of North Carolina, Chapel Hill, NC, USA*

^e *Instituto Superior Técnico, Lisboa, Portugal*

^f *University of Edinburgh, Edinburgh, UK*

^g *Colorado School of Mines, Golden, CO, USA*

^h *Tennessee Technological University, Cookeville, TN, USA*

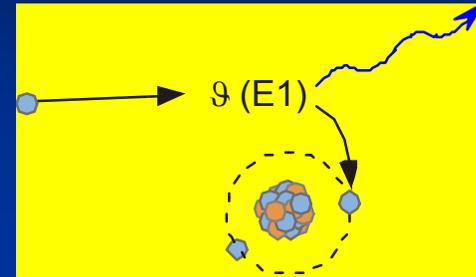
ⁱ *Michigan State University, East Lansing, MI, USA*

^j *Yale University, New Haven, CT, USA*

^k *Rutgers University, New Brunswick, NJ, USA*

Direct capture cross section can be determined by measuring ANC's (or spectroscopic factors) from proton transfer reactions

- Direct capture occurs via an electromagnetic transition at large radii.
- The cross section can be accurately calculated from the Asymptotic Normalization Coefficients (ANC's) with little model dependence.
- The ANC's can be determined by measuring the cross section for peripheral proton transfer reactions.



$$\sigma_{DWBA} \sim |\langle \chi_{\beta} \psi_{\beta} | \mathcal{V} | \chi_{\alpha} \psi_{\alpha} \rangle|^2$$

$$\psi \sim \left(\frac{C}{b}\right) \varphi \quad \text{and} \quad \varphi \xrightarrow{r \gg R_0} b \frac{W}{r}$$

$$\frac{d\sigma}{d\Omega} = \frac{C_{Z+p}}{b_{Z+p}} \frac{C_{17F+p}}{b_{17F+p}} \sigma_{DWBA}$$



Some Remarks on the “ANC Method”

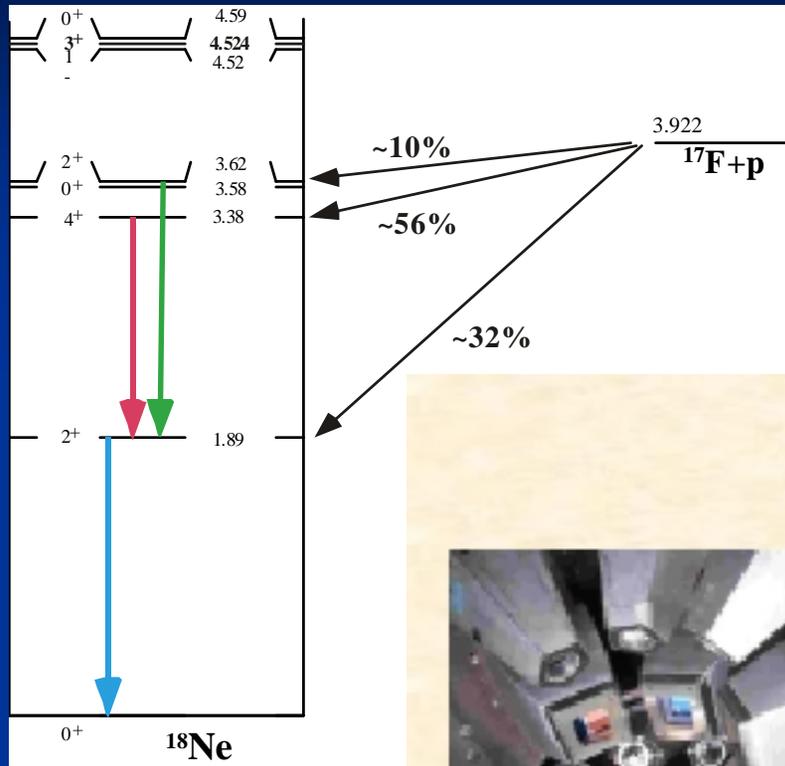
- The idea to use transfer reactions (via spectroscopic factors) to constrain direct capture (DC) has been around since at least the early 1970s.
- Mukhamedzhanov and collaborators contributed the important observation that both the DC and transfer cross section depend mostly asymptotic tail of the bound state (ANC).
- For transfer reactions, the choices of kinematics and reaction can be exploited to reduce theoretical uncertainties. Experimental realities may limit the choices.
- Error analysis must consider model parameters (e.g. optical potentials) as well the reaction mechanism (e.g. 2-step processes, compound-nuclear processes).

Proton transfer reactions are difficult in inverse kinematics (new experimental techniques are required)

- For stable targets the ($^3\text{He},d$) reaction can achieve ~ 15 keV resolution using a magnetic spectrograph.
- Inverse kinematics and low beam intensities (in the case of radioactive ion beams) produce several complications.
- (d,n): gas target? CD_2 target? Neutron detection?
- ($^3\text{He},d$): gas target? Poor kinematics for detecting the deuteron.
- ($^7\text{Li},^6\text{He}$) or ($^{14}\text{N},^{13}\text{C}$)
- The beam-like nucleus can be detected, but energy resolution tends to be poor.
- Gamma-ray tagging can be used for bound excited states.

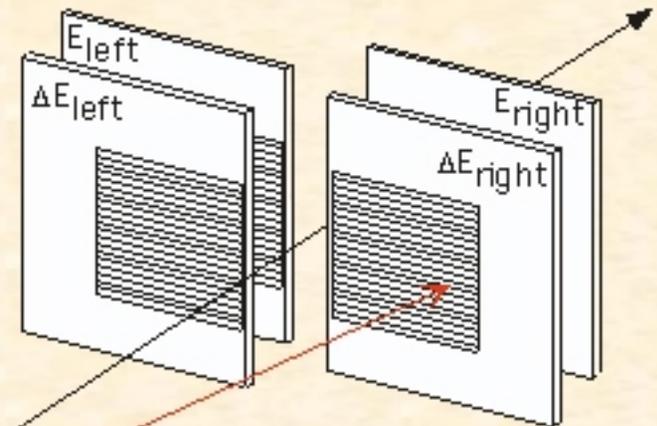
$^{14}\text{N}(^{17}\text{F}, ^{18}\text{Ne}^*)^{13}\text{C}$ at the HRIBF

- The direct capture cross section is dominated by capture to excited states in ^{18}Ne .
- Gamma rays were detected by the CLARION array in coincidence with ^{18}Ne to resolve the states of interest.



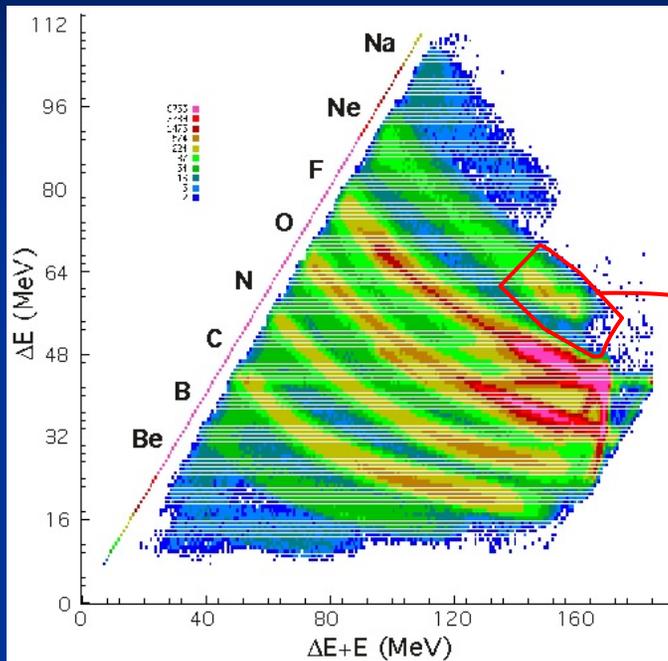
$\text{C}_3\text{N}_6\text{H}_6$ target

^{17}F Beam
(10 MeV/u)

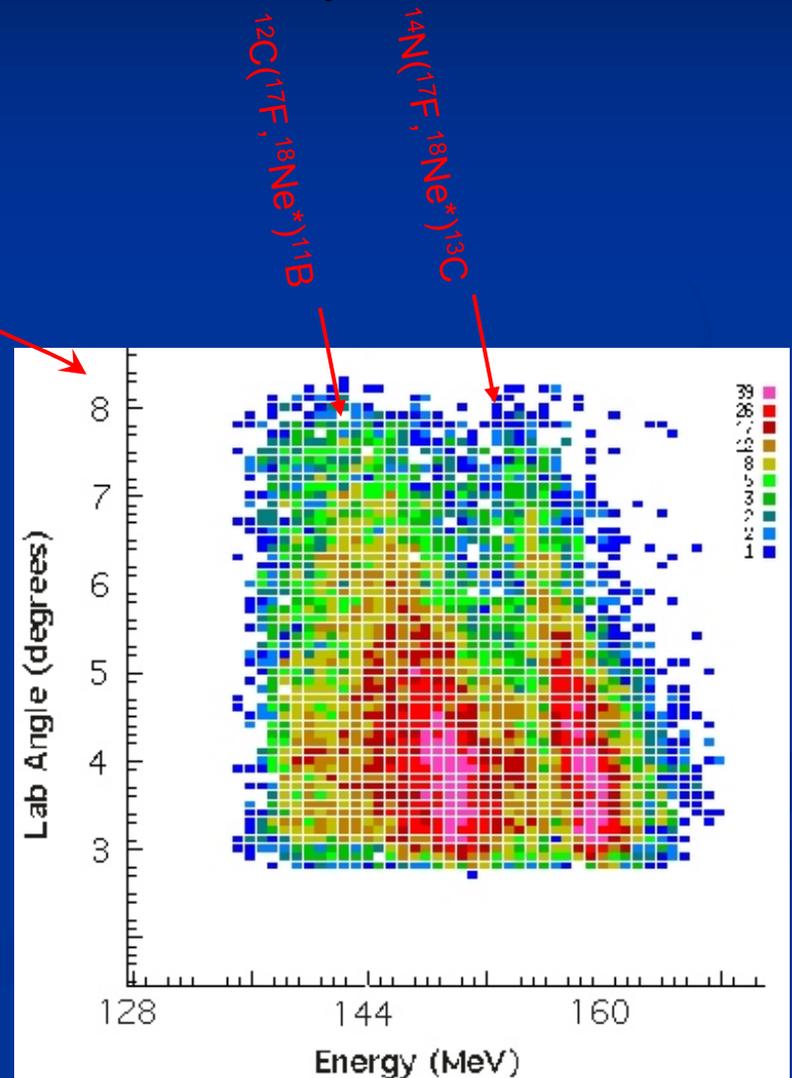


Charged-particle spectra

Particle ID
Summed over whole detector

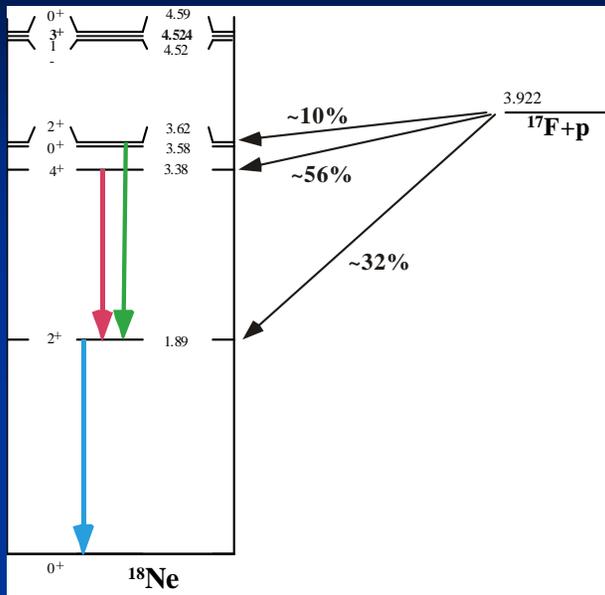


- ^{18}Ne is the strongest neon group, but populated two ways:

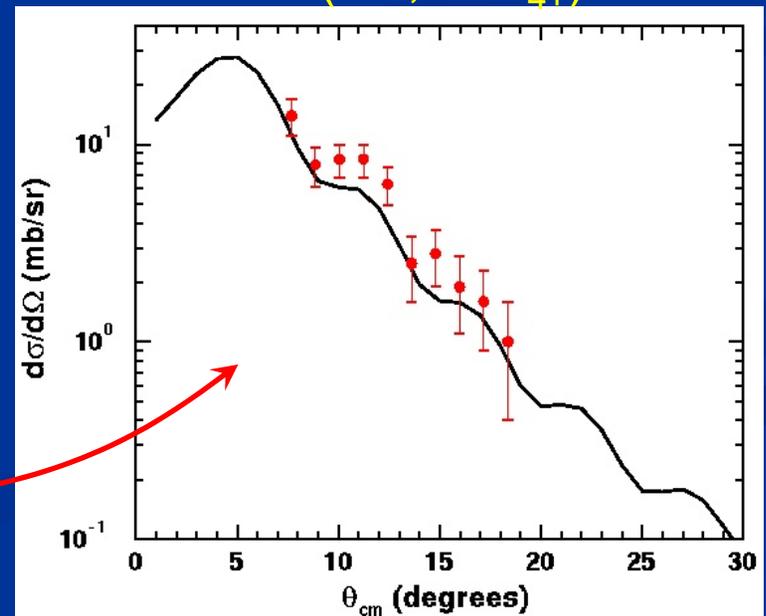
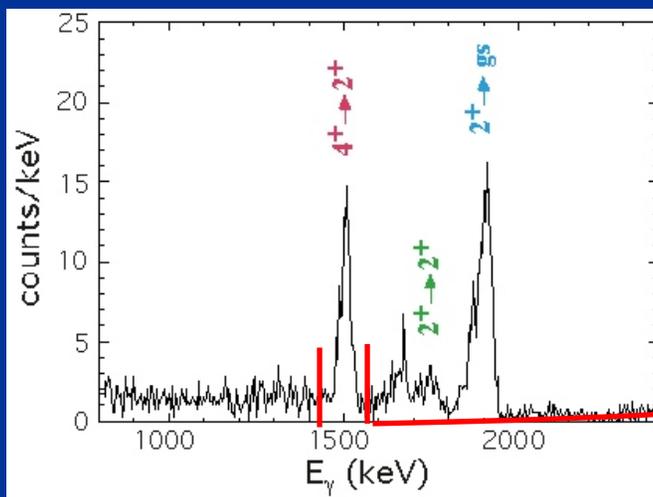


- Good Z separation, but (so far) poor isotopic separation in strip detector
- Charged-particle energy resolution is not good enough (yet) to separate any of the states of interest in ^{18}Ne .

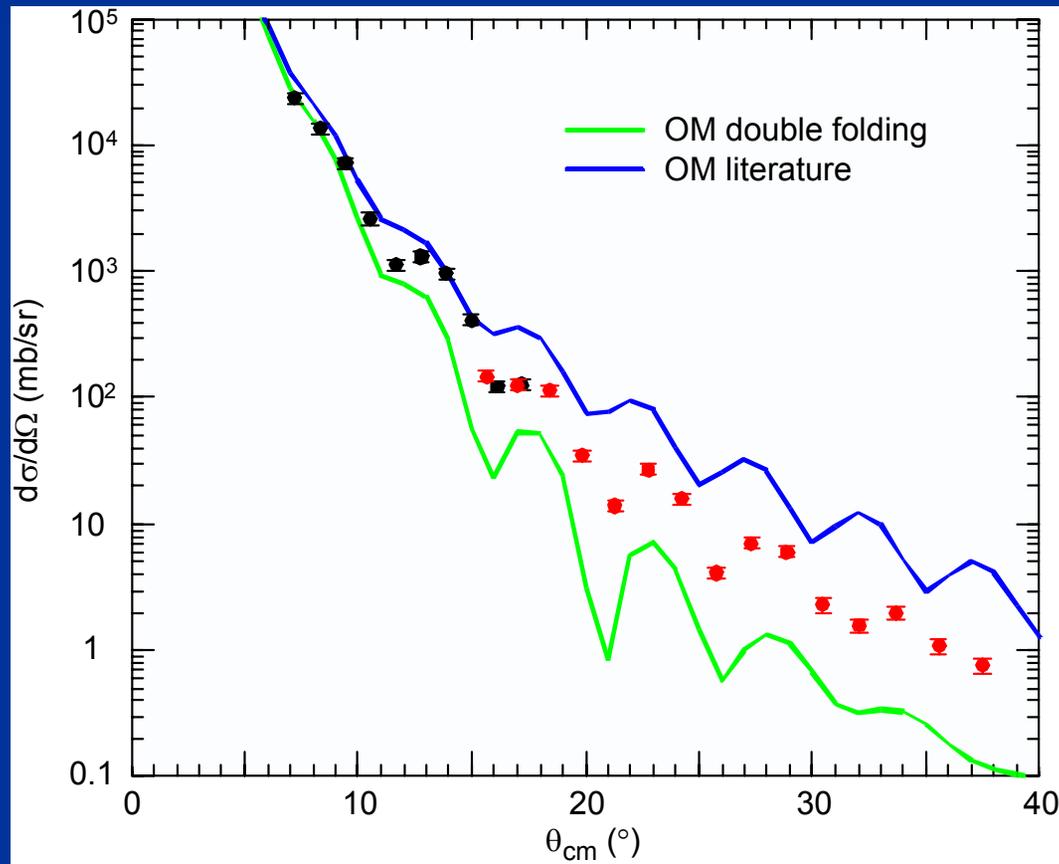
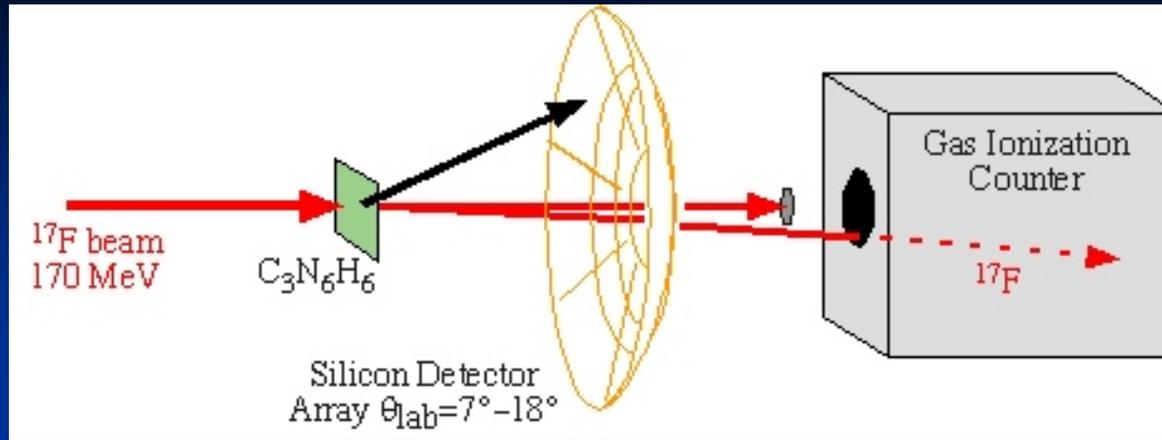
Gammas needed to resolve states



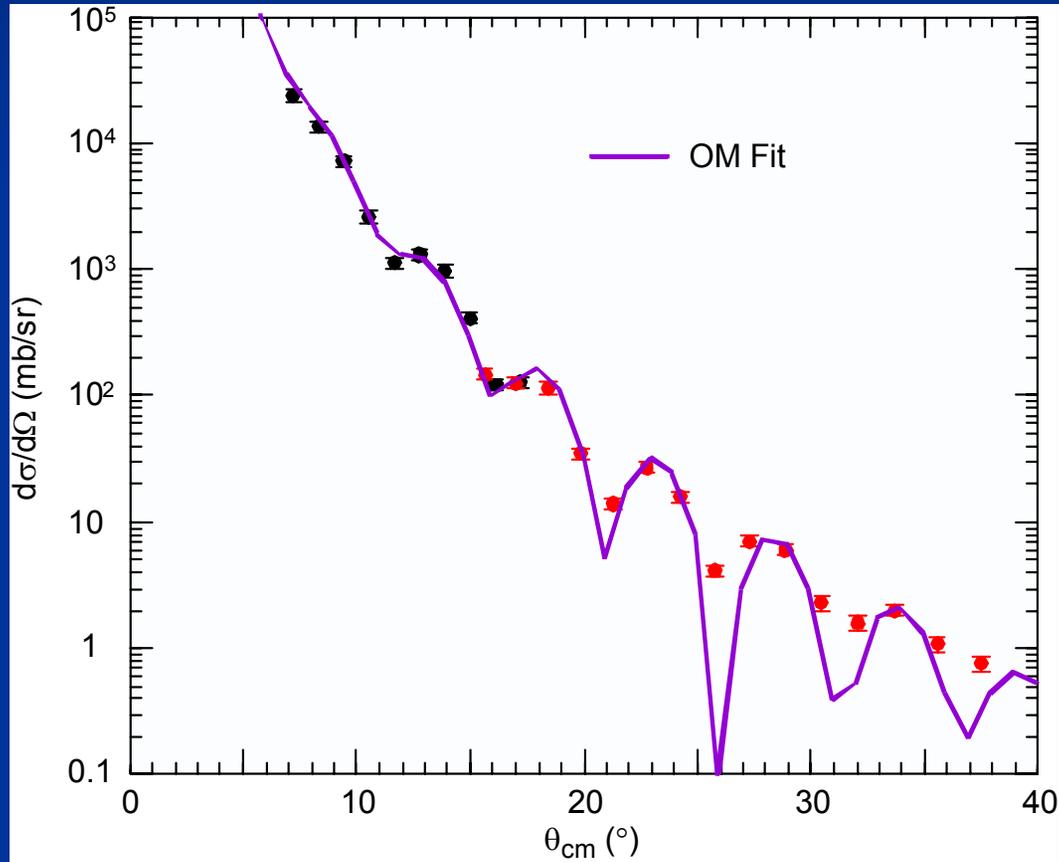
- data analysis in progress
- DWBA calculations with no free parameters
- S factor $\sim 30\%$ higher than shell model calculations
- Expect to accurately determine direct capture cross section within 10 %.



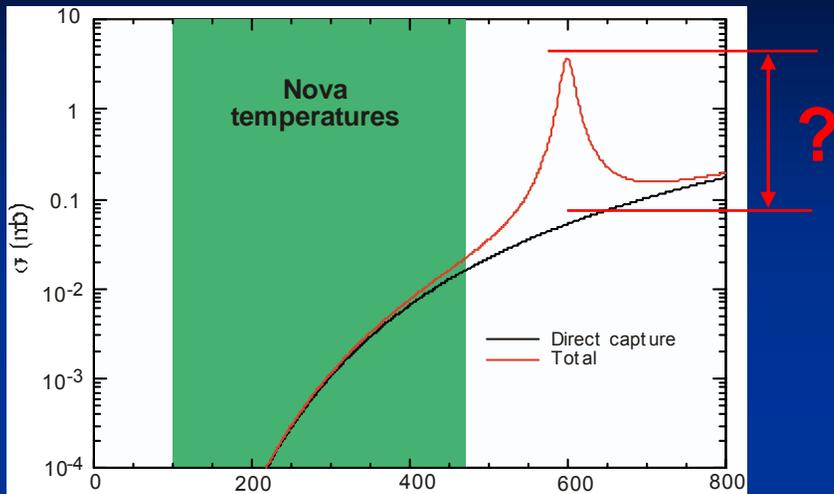
$^{14}\text{N}(^{17}\text{F}, ^{17}\text{F})^{14}\text{N}$ Measurement



These data constrain the Optical Model parameters for the transfer reaction

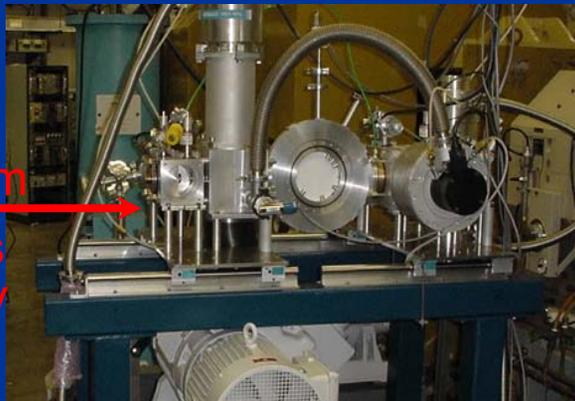


$^{17}\text{F}(p,\gamma)^{18}\text{Ne}$ resonant cross section

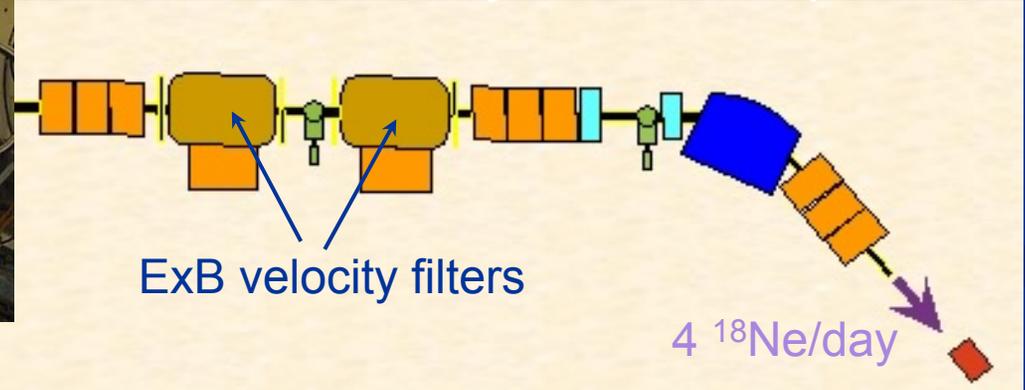


windowless H_2 gas target
10 mg/cm²

- Amplitude of 3^+ resonance is uncertain.
- Dominates the reaction rate at higher temperatures.
- We hope to measure the $^{17}\text{F}(p,g)^{18}\text{Ne}$ resonant cross section using a H_2 gas target and the DRS.



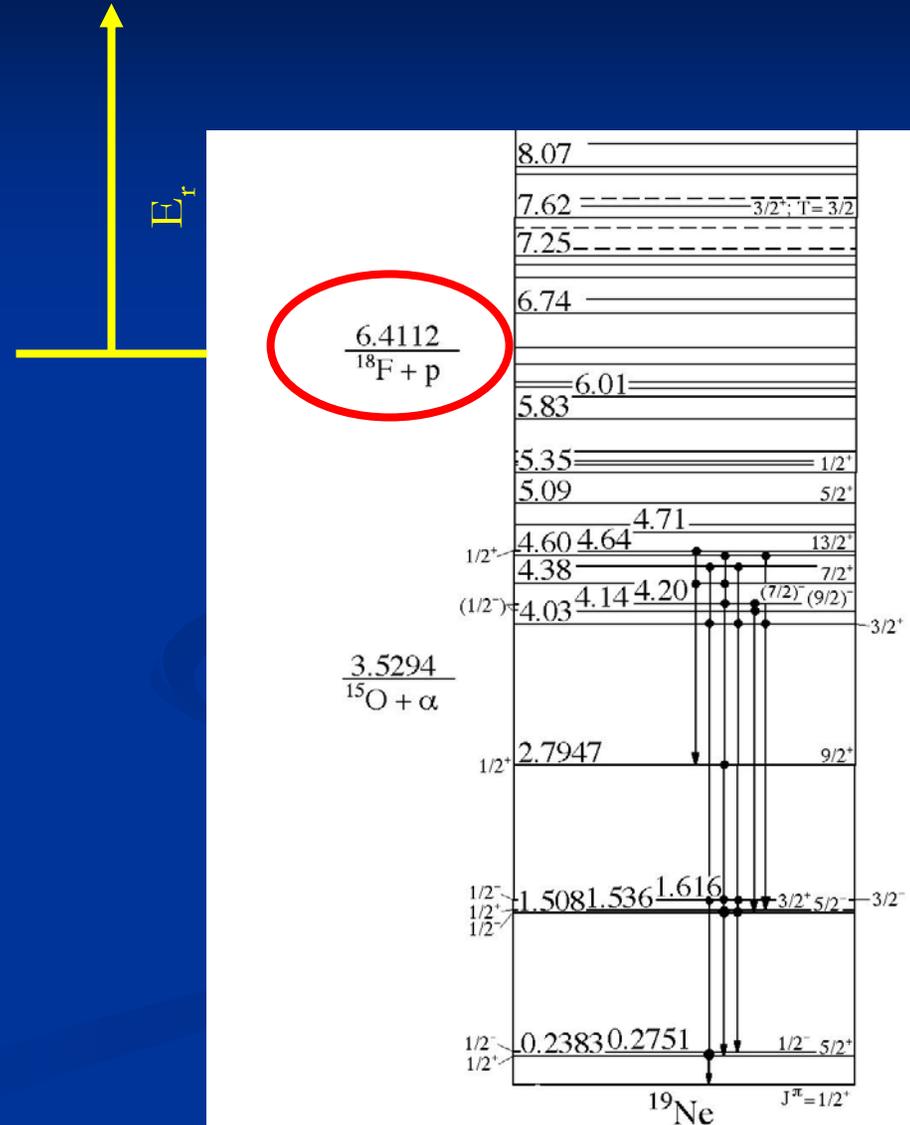
The Daresbury Recoil Separator



- ^{18}Ne from $^1\text{H}(^{17}\text{F},^{18}\text{Ne})$ reaction detected by gas ionization counter.

$^{18}\text{F}(p,\alpha)^{15}\text{O}$ and $^{18}\text{F}(p,\gamma)^{19}\text{Ne}$

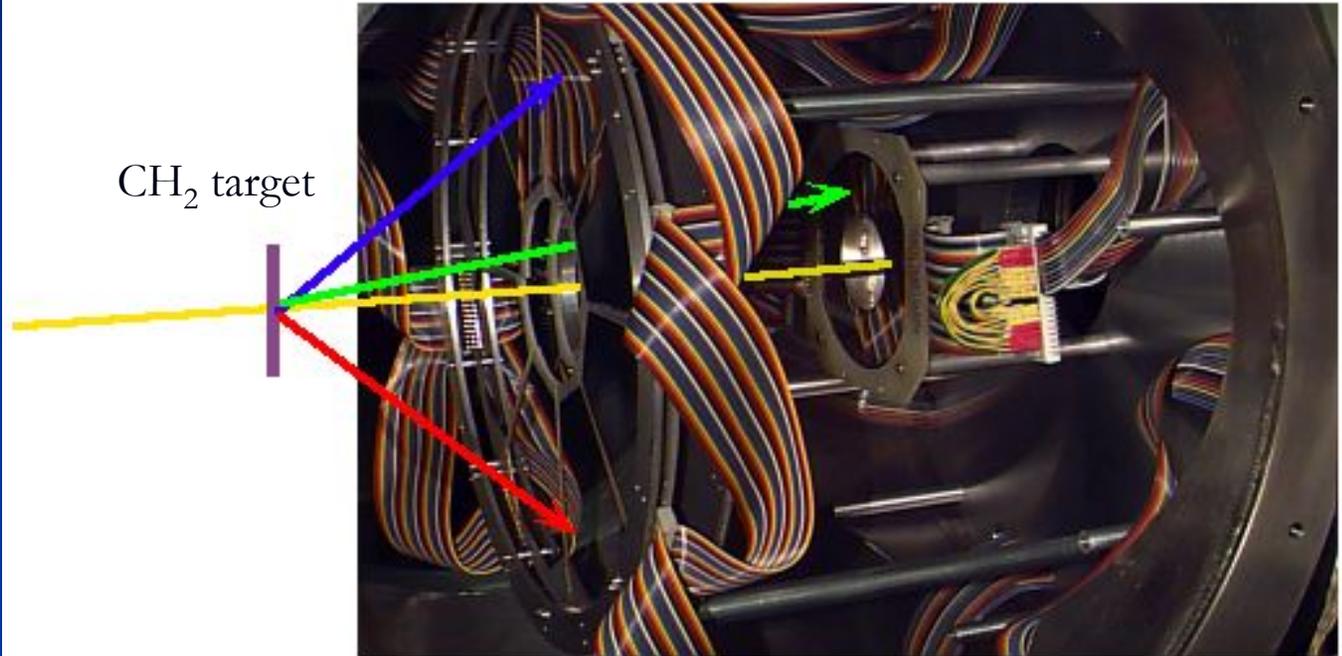
- Several resonances may be important for nova temperatures
- $^{18}\text{F}(p,\alpha)$ can be measured directly, but not over the entire energy range needed for novae.
- Transfer reactions and mirror symmetry can also be used.



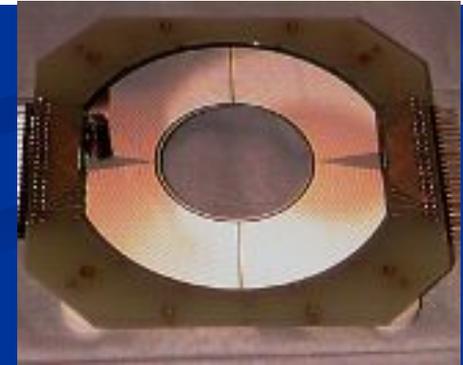
Experimental Approach

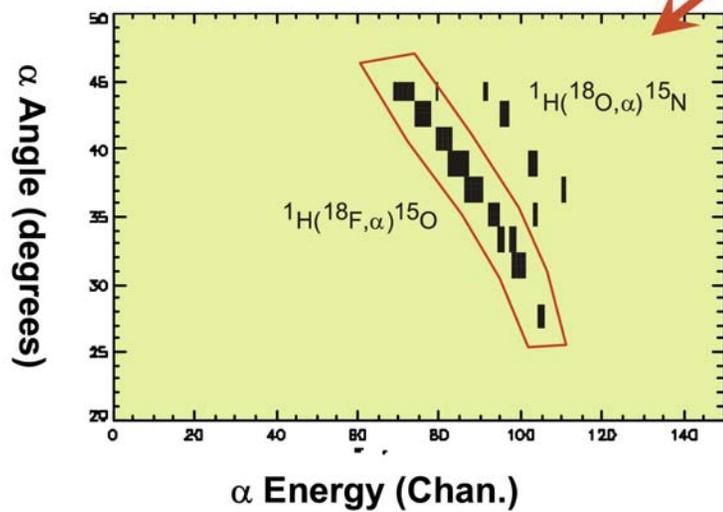
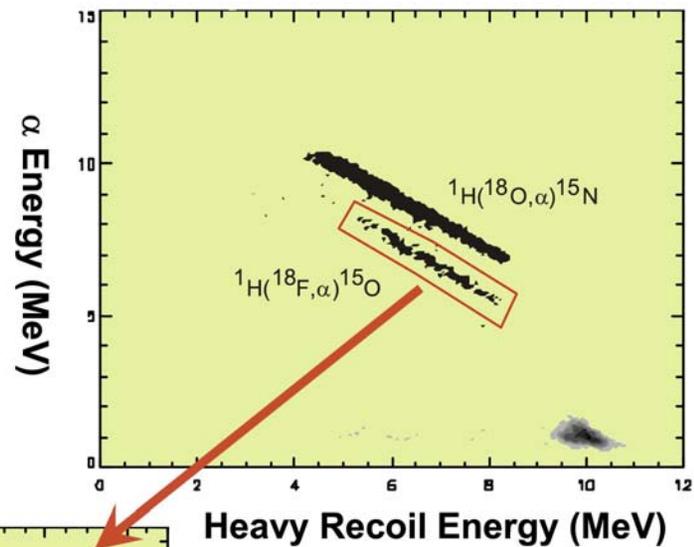
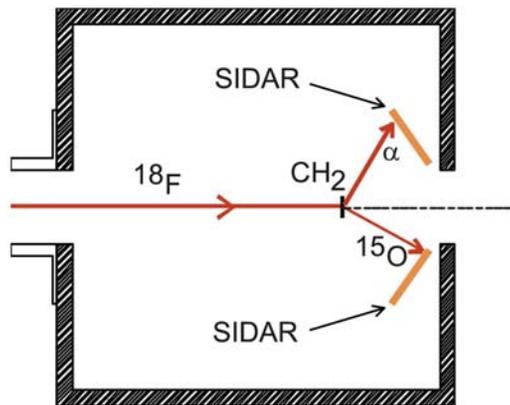
$\sim 5 \times 10^5 \text{ }^{18}\text{F}/\text{sec}$

CH₂ target

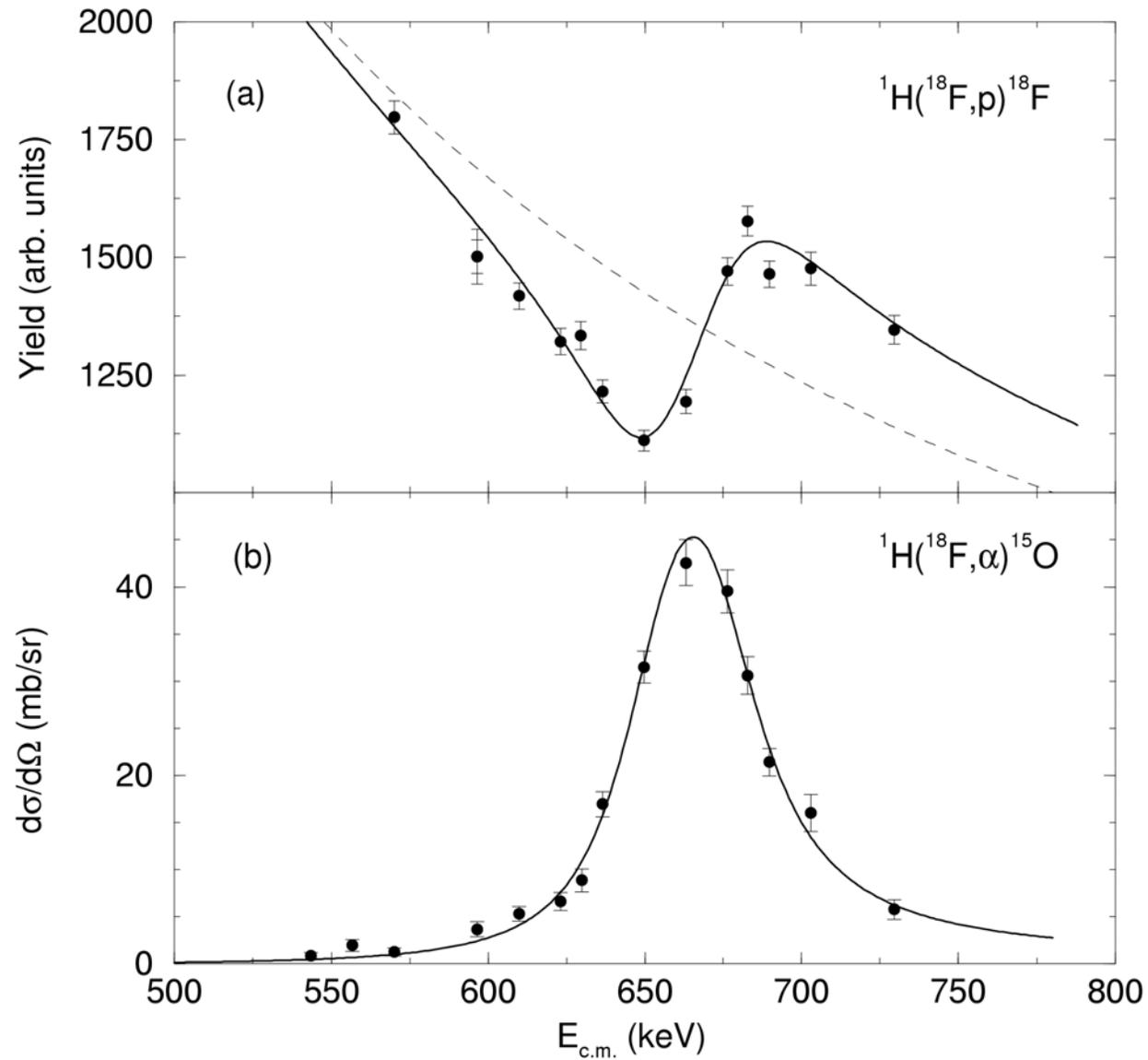


Si Strip Detectors

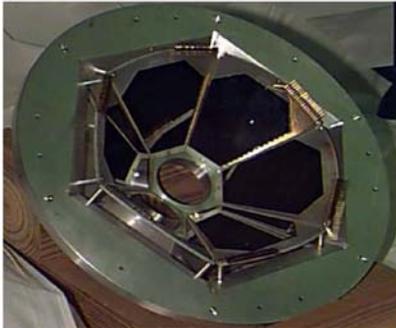
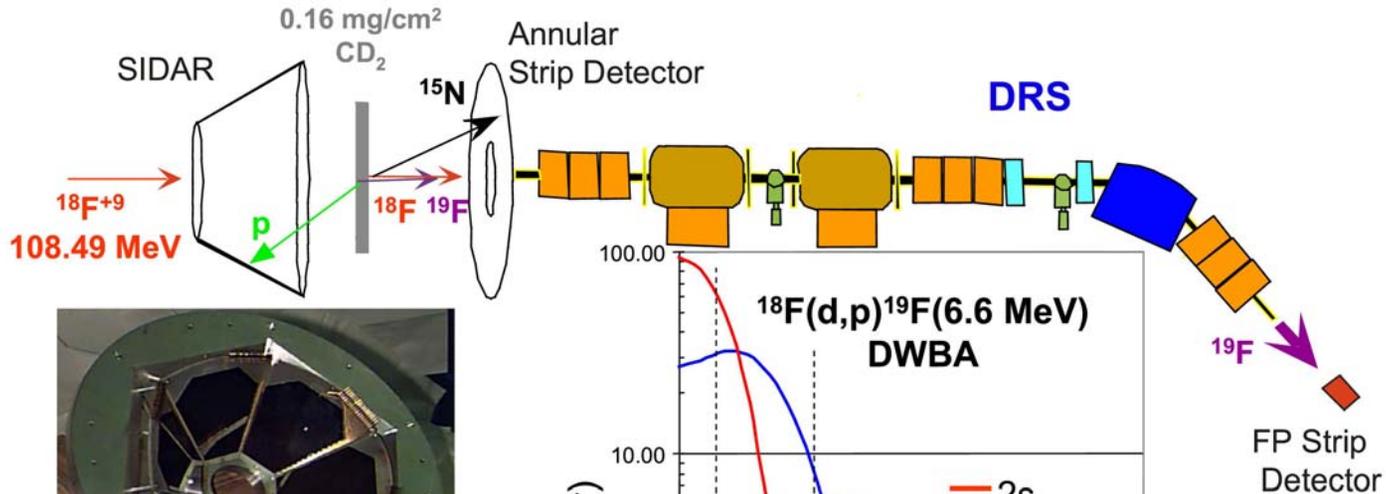




Results



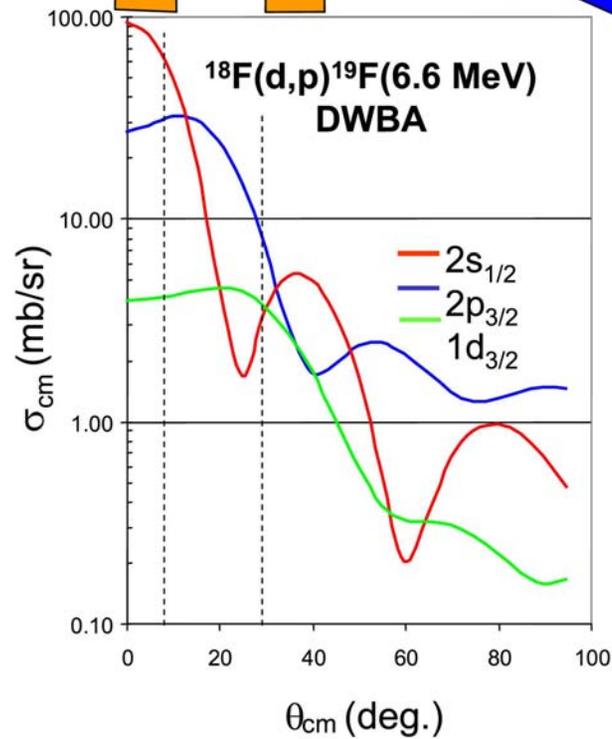
${}^2\text{H}({}^{18}\text{F},\text{p}){}^{19}\text{F}$ at the HRIBF (6 MeV/u)



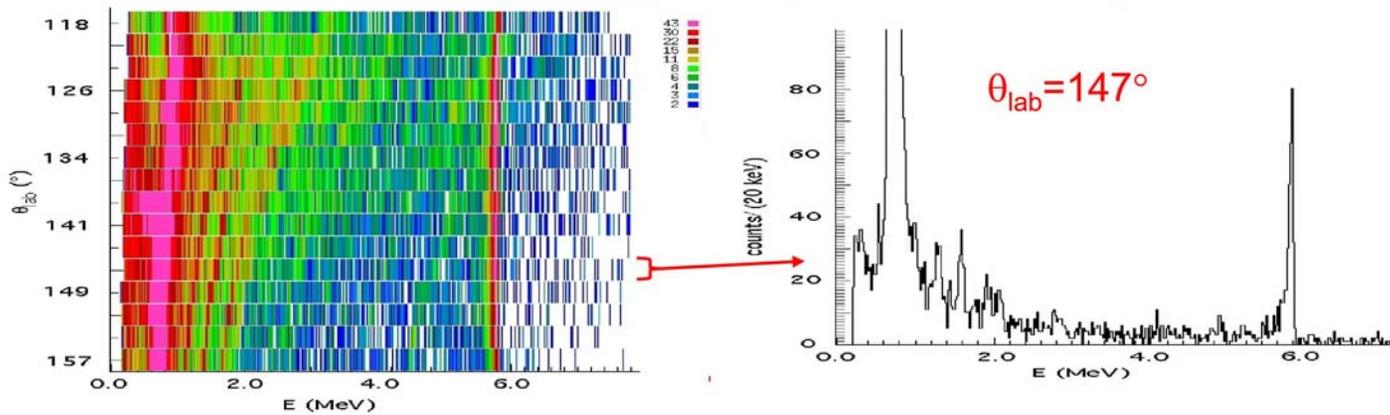
$$\theta_{\text{lab}} \approx 116^\circ - 160^\circ$$

$$\theta_{\text{cm}} \approx 7^\circ - 29^\circ$$

~ 3 days of data with
 5×10^5 ${}^{18}\text{F}/\text{s}$ on target.

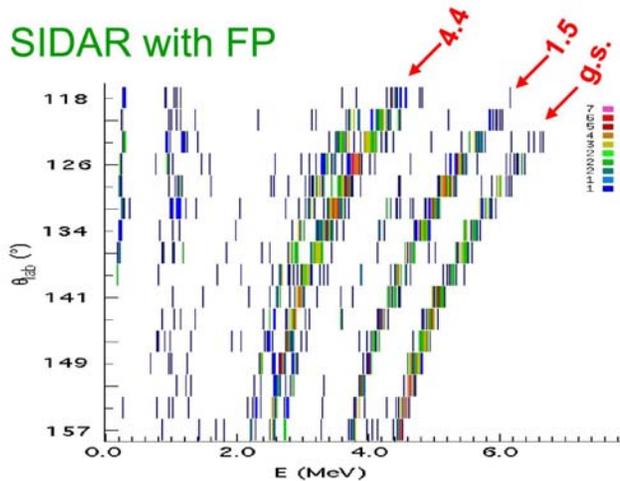


SIDAR Singles Spectra (One Detector)

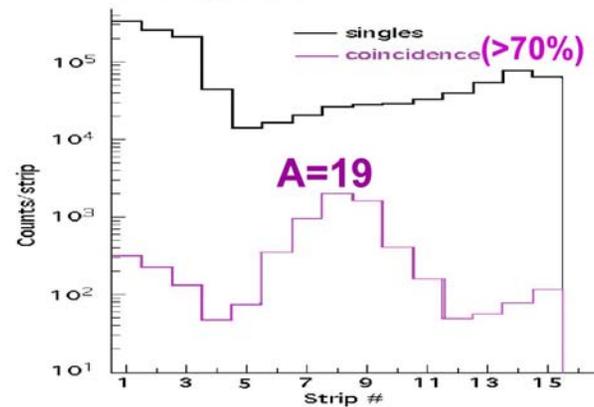


SIDAR-FP Coincidences

SIDAR with FP

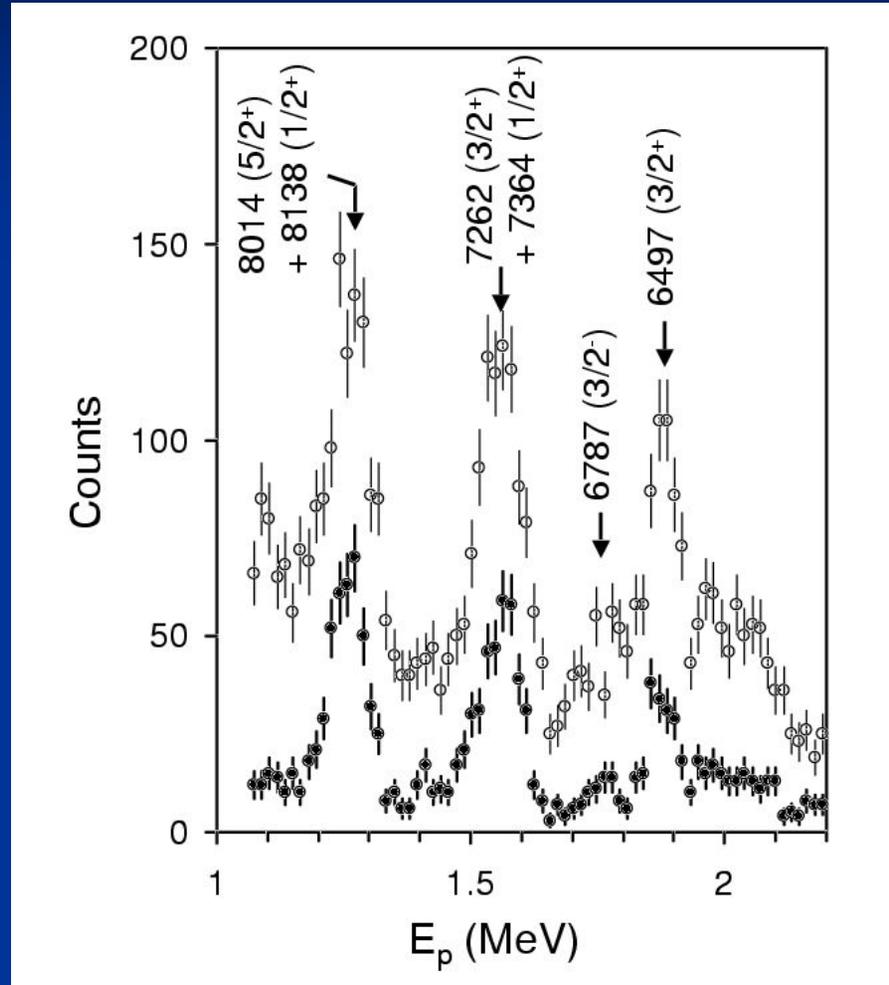


FP Position



Results

$^{18}\text{F}(d,p)^{19}\text{F}$
proton spectrum
 $\theta_{\text{lab}} = 147^\circ$

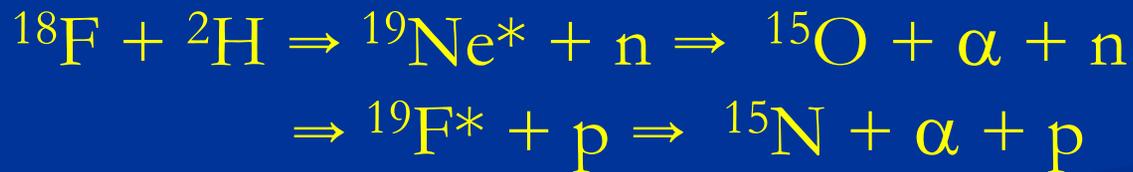


Proton Transfer on ^{18}F

Appeared to be difficult...

but the ^{19}Ne states of interest break up into $^{15}\text{O} + \alpha$ which provides a unique signature.

Our new approach to (d,n) and (d,p) :

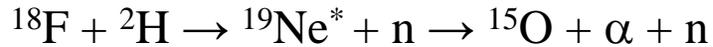


without detecting the n or p.

- The ^{15}O and α are detected with position-sensitive Si strip detectors.
- The relative energy can thus be reconstructed.
- This approach is less sensitive to target thickness (720 $\mu\text{g}/\text{cm}^2$ was used).
- Work of my student: Remi Adekola

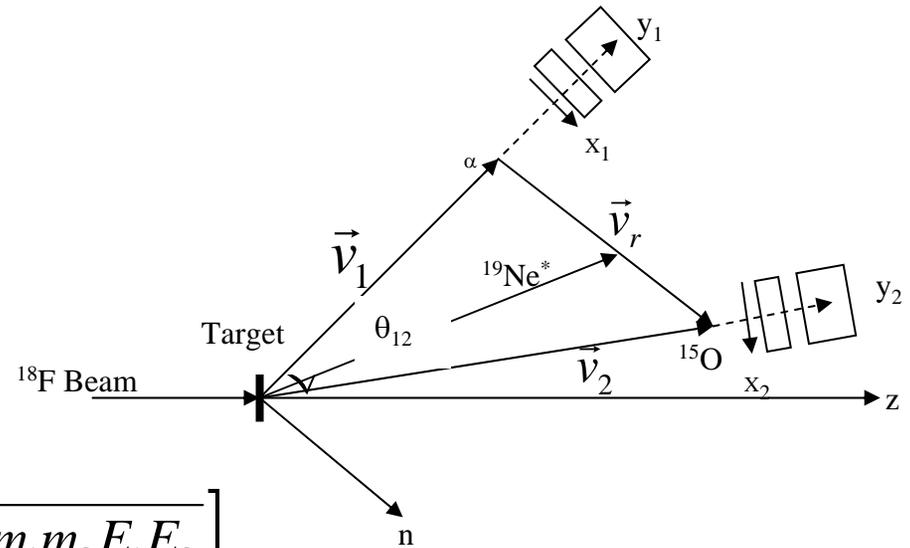
Reconstructing the Relative Energy

Reaction:



relative energy of the state

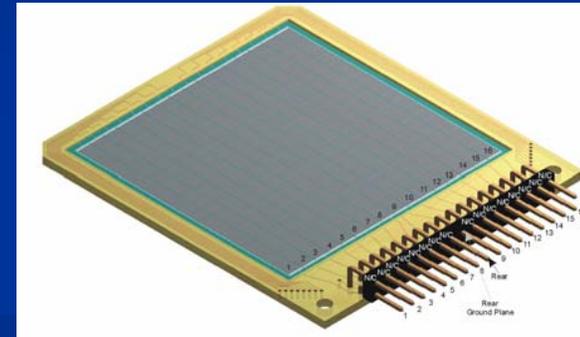
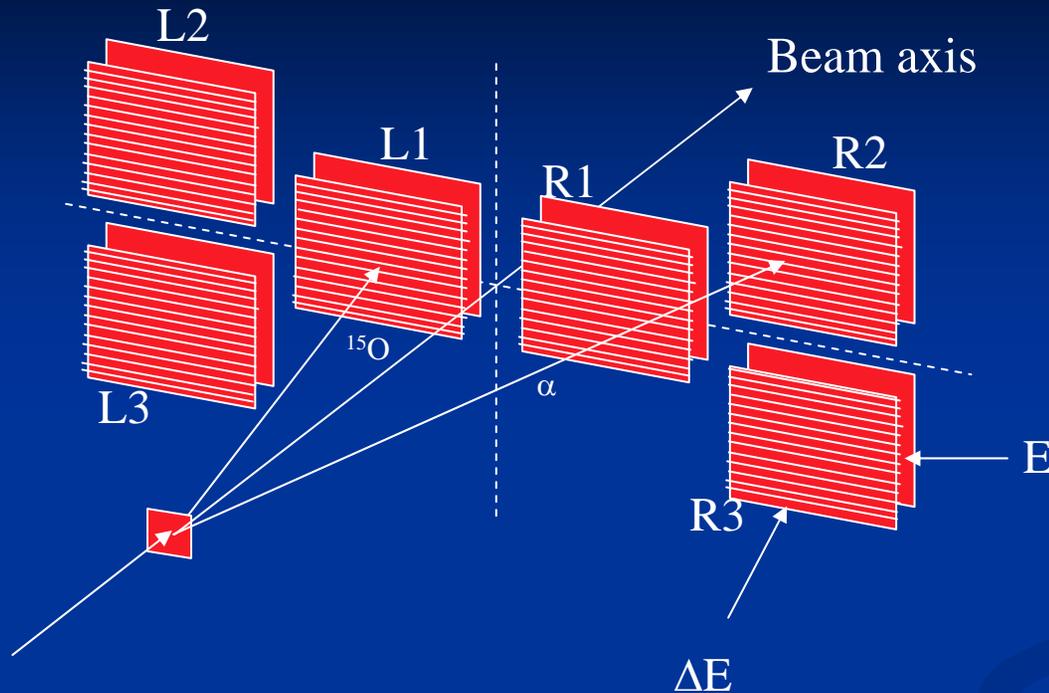
$$E_{rel} = \frac{1}{m_1 + m_2} \left[m_1 E_2 + m_2 E_1 - 2 \cos \theta_{12} \sqrt{m_1 m_2 E_1 E_2} \right]$$



$$E_x = E_{th} + E_{rel}$$

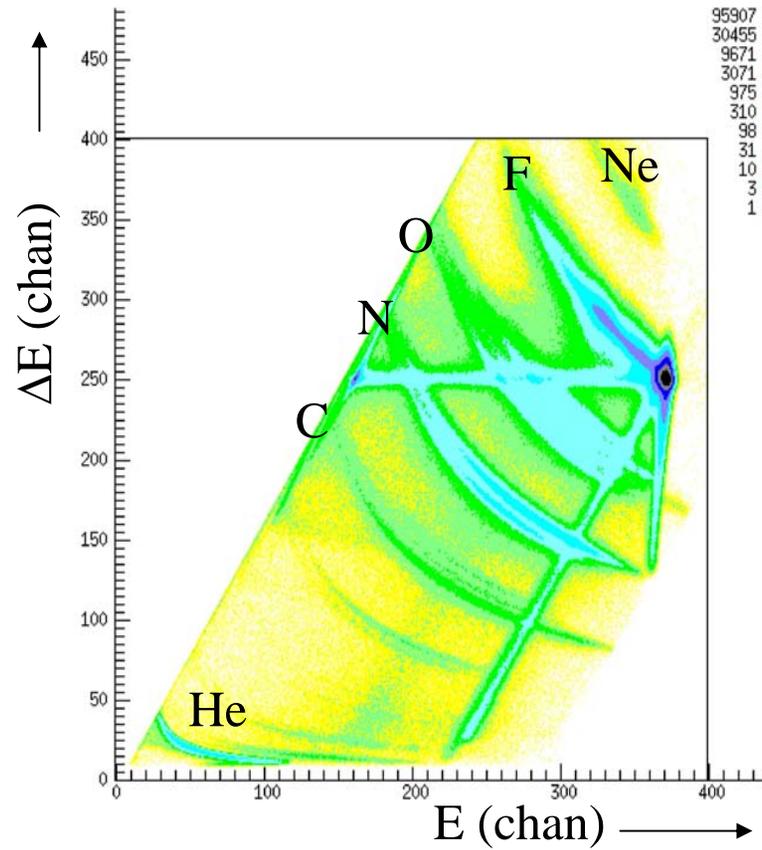
$$Q = E_1 + E_2 + E_3 - E_A$$

Detector configuration



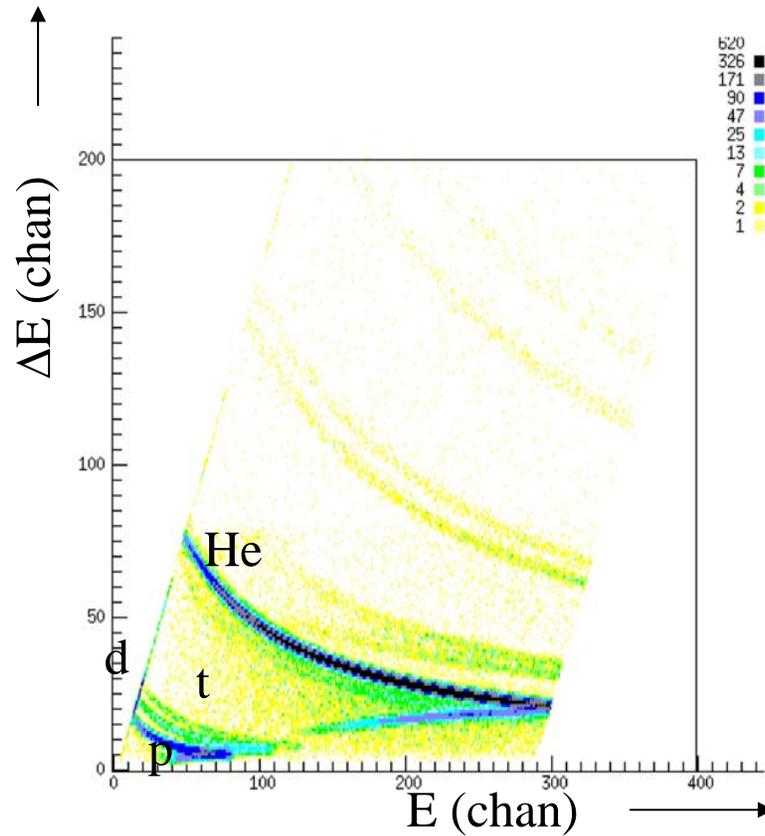
- Each telescope is 5 cm x 5 cm and located ~45 cm downstream from the target.
- Inner ΔE s are 65 μm ; others are 140 μm ; E detectors are 1 mm.

Inner Telescope

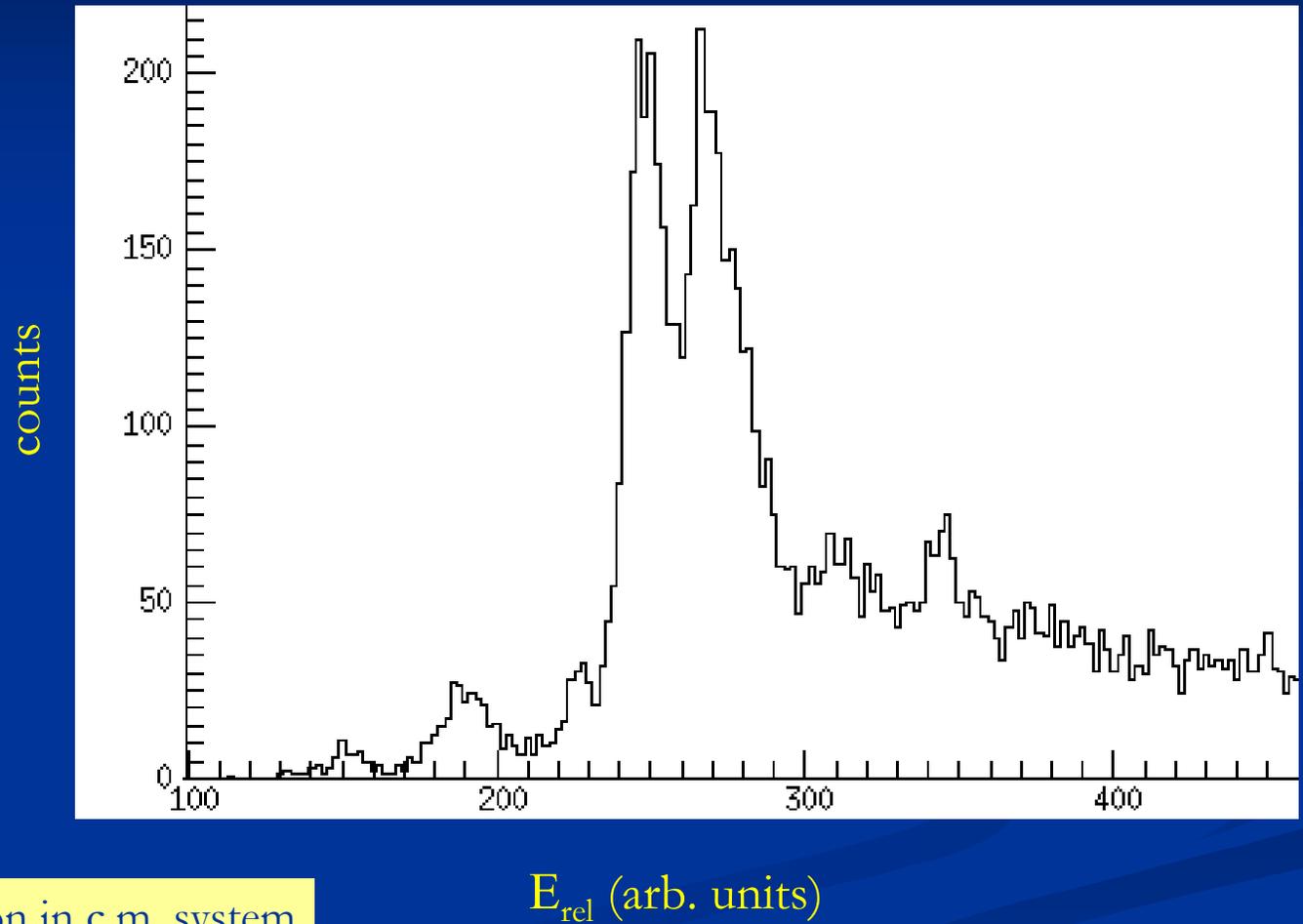


Particles identification histogram

Outer Telescope



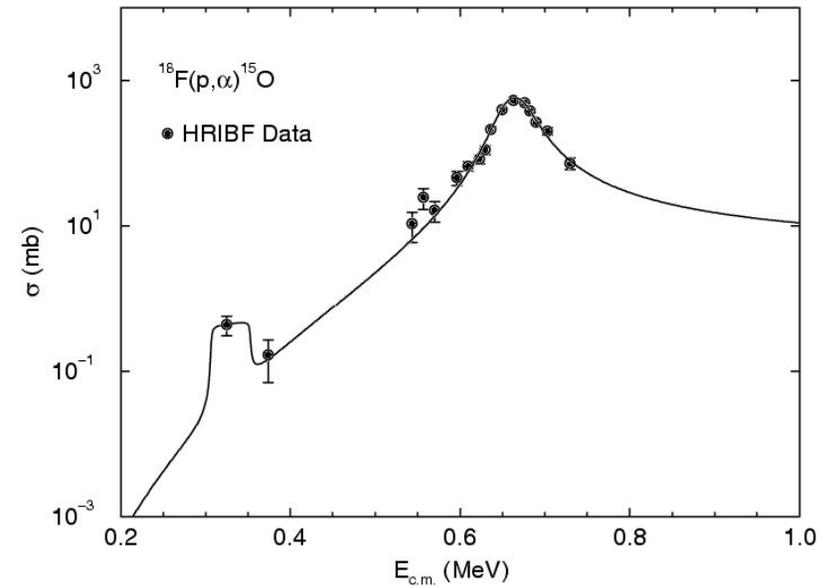
Preliminary Relative Energy Spectrum



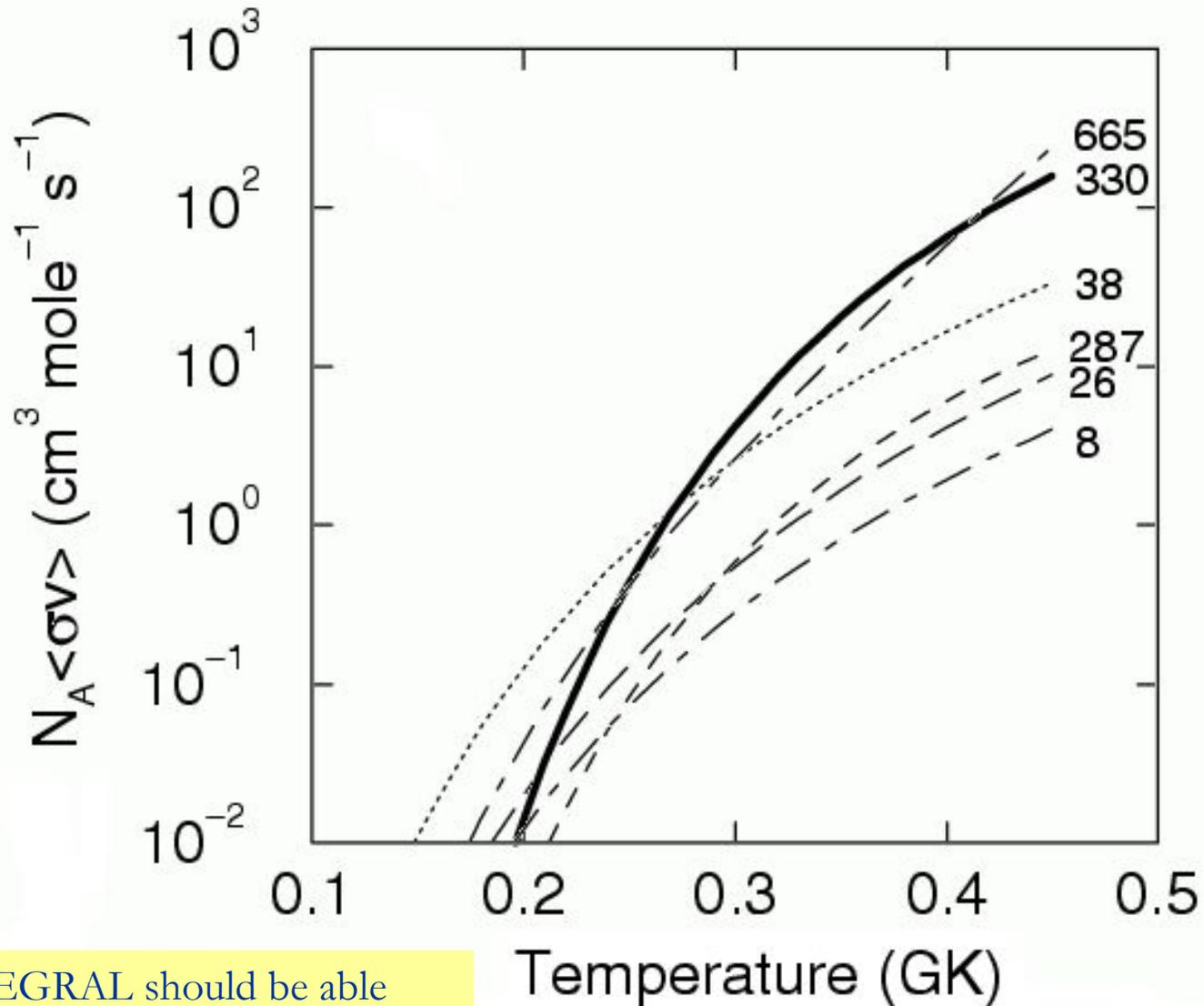
Energy resolution in c.m. system
is ~ 70 keV.

Our Present Understanding

E_r (keV)	J^π	Γ_p (keV)
8	$3/2^+$	4×10^{-37}
26	$1/2^-$	3×10^{-20}
38	$3/2^+$	2×10^{-14}
287	$5/2^+$	4×10^{-5}
330	$3/2^-$	$2.2(0.7) \times 10^{-3}$
665	$3/2^+$	$15.2(1.0)$

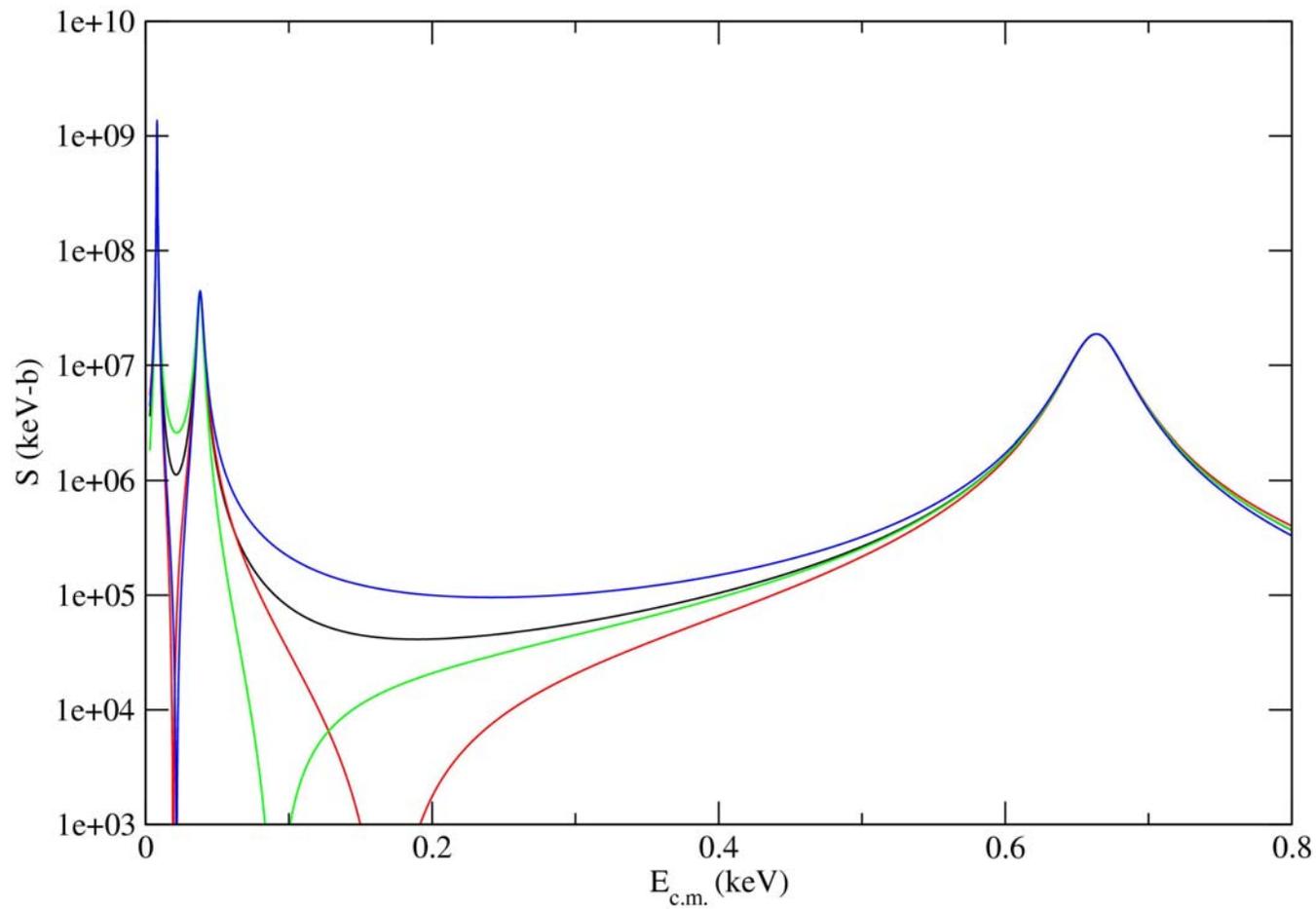


Reaction Rate



Note: SPI/INTEGRAL should be able to see 511-keV photons following a nova outburst provided it is with ~ 5 kpc of earth.

Interfering $3/2^+$ Resonances



For the Future:

- Complete analysis of proton transfer data.
- Measure $^{18}\text{F}(p,\alpha)^{15}\text{O}$ at lower energies?
- Measure spectroscopy with $^{17}\text{O}(^3\text{He},n)^{19}\text{Ne}$.

Thanks: D.W. Bardayan, J.C. Batchelder, J.C. Blackmon, W. Bradfield-Smith, A.E. Champagne, J.A. Cizewski, T. Davinson, U. Greife, C.J. Gross, M. Hornish, C. Iliadis, C.C. Jewett, B.A. Johnson, R. Kozub, C.S. Lee, R. Lewis, R.J. Livesay, Z. Ma, T.N. Massey, C. Matei, B.H. Moazen, C.D. Nesaraja, P.D. Parker, L. Sahin, J.P. Scott, D. Shapira, N. Shu, M.S. Smith, J. Thomas, D.W. Visser, A. Voinov, P.J. Woods

